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# **Dynamics and Bifurcations in Fractional Lozi Map**

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**Received:** 6 February 2023; **Revised:** 27 October 2023; **Accepted:** 27 October 2023

**Abstract:** We propose a fractional version of the Lozi map. The Lozi map is quite similar to the Henon map except that the square term of the Henon map is replaced by a linear term in the Lozi map. This map is known to be analytically tractable. The fractional version of the Lozi map displays rich dynamical behaviour. In this work, we carry out a timeseries analysis of the map and present the bifurcation diagrams along with attractors for different parametric values of the map for various values of fractional order parameter  $(\alpha)$ .

*Keywords***:** Lozi map, henon map, fractional calculus, chaotic maps, coupled map lattice

**MSC:** 37E99, 26A33

#### **1. Introduction**

Nonlinear maps manifest various phenomena that occur in flows as well. We can observe chaos even in onedimensional nonlinear maps while the chaos is not possible in one or two-dimensional differential equations. In some cases, it is possible to derive maps from flows using the Poincaré map. Nonetheless, higher dimensional maps show dynamically rich behaviour.

One of the earliest and most studied chaotic maps is the Henon map [1]. Later on, a more analytically tractable version of the map known as the Lozi map was introduced and has been studied in great detail [2]. The Lozi map is piecewise linear. Piecewise linear maps have been studied in great detail recently due to their application in power electronics [3]. Bifurcations and the onset of chaos in nonlinear maps have been an object of immense interest to both mathematicians and physicists. We study a fractional version of the Lozi map in this work. [O](#page-6-0)f late, fractional versions of nonlinear chaotic maps have attracted much attention [4–9]. Pioneering work has been done by Miller an[d R](#page-6-1)oss [10], Atici and Eloe [11] and notable others [12, 13]. The complex nature of fractional nonlinear models has also been investigated in [14]. [A](#page-6-2) mechanical system can also be modelled using a fractional nonlinear map like the chaotic water wheel model in [15]. Another notable work investigating nonlinear maps of fractional order is [16]. Due to the possible applications of these systems in varied engineering and co[ntr](#page-6-3)[ol](#page-6-4) processes, they are studied with interest. A nonlinear m[ap o](#page-6-5)r difference equat[ion](#page-6-6) as it is studied often [do](#page-6-7)[es n](#page-6-8)ot have any long-term memory. This is because the current state of the system depen[ds u](#page-6-9)pon the previous state of the system only. If we want to model a system with memory, the evolution should explicitly dep[end](#page-6-10) upon all the previous states of the system. This explicit dependence of the [cu](#page-6-11)rrent state on the previous states introduces

DOI: https://doi.org/10.37256/cm.5420242423

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long-term memory in the system. All definitions of the fractional difference equations include this explicit dependence on the previous states. Thus fractional difference equations are suitable mathematical models for systems having long-term memory. This makes the study of fractional nonlinear systems very interesting and useful from the point of view of many engineering and scientific applications. Chaos of the fractional difference maps has been the focus of a recent study by Xiao et al. [17]. Notable work has been done by Gejji [18], Baleanu [19], Atici [20], Edelman [21], etc on fractional difference equations.

Here the analysis done is mainly numerical. However, the extension of the theory of Lyapunov functions to fractional maps will be of great interest. It has been carried out for fractional differential equations [22].

We also [no](#page-7-0)te that unstable periodic orbits play an imp[orta](#page-7-1)nt role in t[he d](#page-7-2)ynamics [of](#page-7-3) Lozi map. [How](#page-7-4)ever, for fractional order maps, we have results only for fixed point and period-2 orbit [23]. The Lozi map is similar to the Henon map. The difference is the presence of the piecewise linear term in place of the quadratic term in the Henon map. Lozi introduced a two-dimensional map in 1978 [24]. It resembles the Henon map [1] with the quadratic t[erm](#page-7-5) in the Henon map replaced with a linear relation. The integer order Lozi map is given as,

$$
x(t) = 1 + y(t-1) - a|x(t-1)|,
$$
\n(1)

$$
y(t) = \beta x(t-1),\tag{2}
$$

where *a* and  $\beta$  are the map parameters that could be varied. The Lozi map is a modified version of the Henon map where the quadratic term in the Henon map is replaced by a linear term. For  $\beta = 0.05$  and  $\beta = 0.2$ , we have shown a bifurcation diagram in Figure 1. Some studies on fractional versions of the Henon map (although with slightly different definitions) and Henon-Lozi map have been carried out previously [25, 26].



**Figure 1.** Bifurcation diagram for (a)  $\beta = 0.05$  and (b)  $\beta = 0.2$ 

We expect the dynamics to be different for the Lozi map than the Henon map. The dynamical behaviour of integer order Lozi map has been investigated extensively. The fractional version of the Lozi map displays rich dynamical behaviour like abrupt order-to-order or order-to-chaos transitions, infinite sets of period-3 cycles embedded in the sea of chaos [27]. We also note that analytical studies are possible in the Lozi map due to its piecewise linear nature. For example, Misiurewicz proved the conjecture of Lozi about the existence of a chaotic attractor for the Lozi map [28].

We define the Lozi map and show its bifurcation diagram for some parameters so that we can compare these dynamics to a fractional version of the Lozi map. This would help us to understand the impact of long-term memory on the dynamics of these systems.

### **2. Fractional Lozi map**

We use the definition of fractional difference operator by Deshpande and Daftardar-Gejji [18]. According to their definition, the evolution of a 1-d fractional nonlinear map is given as,

$$
x(t) = x(0) + \frac{1}{\Gamma(\alpha)} \sum_{j=1}^{t} \frac{\Gamma(t - j + \alpha)}{\Gamma(t - j + 1)} [f(x(j-1)) - x(j-1)],
$$
\n(3)

where  $f(x)$  is a nonlinear map and  $\alpha$  is the fractional order parameter which can have any real number values between 0 and 1. We extend this definition to 2-d by replacing *x* with vector (*x, y*). A similar definition can be observed in [29]. The fractional version of the Lozi map is then given as,

$$
x(t) = x(0) + \frac{1}{\Gamma(\alpha)} \sum_{j=1}^{t} \frac{\Gamma(t - j + \alpha)}{\Gamma(t - j + 1)} [1 + y(j - 1) - a] x(j - 1)] - x(j - 1)],
$$
\n(4)

<span id="page-2-1"></span><span id="page-2-0"></span>
$$
y(t) = y(0) + \frac{1}{\Gamma(\alpha)} \sum_{j=1}^{t} \frac{\Gamma(t - j + \alpha)}{\Gamma(t - j + 1)} [\beta x(j - 1) - y(j - 1)].
$$
\n(5)

As is evident from the above definition, the evolution of a current state explicitly depends upon all the previous states of the system. This is the reason the fractional nonlinear system is said to have memory. There are two map parameters, namely *a* and β.

We evolve the system defined by eq. (4) and (5) in time for various values of fractional order parameter  $\alpha$  and map parameter β. We discard 10<sup>4</sup> initial iterates of the system and use the last 10 values of variable *x* to plot the bifurcation diagram as a function of map parameter *a*. Even if we discard the first 10<sup>5</sup> transients, the bifurcation diagram is unchanged. We observe that, for  $\alpha = 0.4$ , the system tends to go to infinity above  $\beta = 0.1$ . For  $\alpha = 0.6$ , the system goes to infinity above  $\beta = 0.26$  [a](#page-2-1)nd for  $\alpha = 0.8$ , it goes to [in](#page-2-0)finity above  $\beta = 0.54$ . Thus we can explore a larger range of  $\beta$  for a larger value of  $\alpha$ . We take  $x = y = 0$  as the initial conditions. For  $\beta = 0.05$ , we observe a period-1, followed by period-2 and then period-4. Period-4 attractor develops into 4-band attractor which continuously merges and leads to 2-band attractor. The 2-bands continuously grow and merge to form a single attractor. The same dynamics is seen for larger  $\beta$  except that the fixed point is not observed. We observe that for larger dissipation i.e. for a large value of  $\beta$  the size of the chaotic attractor shrinks. The bifurcation diagrams for  $\beta = 0.05, 0.08$  and 0.1 are presented in Figure 2. We can infer that for a larger value of fractional order parameter  $\alpha$ , the system is more stable and doesn't go to infinity for a wider range of β values. For the lesser value of  $\alpha$ , the system tends to go to infinity for lesser values of  $\beta$ . Thus for the higher values of the fractional order parame[te](#page-3-0)r  $\alpha$ , the chaos is much robust. Increase in fractional order parameter  $\alpha$  generally makes the system much more stable.

<span id="page-3-0"></span>

**Figure 2.** Bifurcation diagrams for fractional Lozi map for various values of fractional-order parameter  $α$  and map parameter  $β$ 

The attractors can be visualized by the time-series plots. We iterate the system for  $1.5 \times 10^4$  time steps and plot the variable values as a function of time. We observe that, as the map parameter *a* is increased, the system goes through a transition from a periodic regime to a chaotic regime in the manner described above. The period-1, period-2, period-4, 4-band, 2-band and single chaotic attractor can be observed from the time series given in Figure 3.



**Figure 3.** Time series plots for the fractional Lozi map showing (a) period-1 (b) period-2 (c) period-4 (d) 4-band (e) 2-band and (f) single band

We study the attractors for the system for fractional order parameter  $\alpha = 0.4$  and map parameter  $\beta = 0.05$  in further detail. The attractors shown in the figures below clearly show the transition as the value of *a* is changed from 1*.*29 for period-4 bifurcation to a single band for  $a = 1.39$ . We have shown period-4, 4-band, 2-band, and single band attractor in Figure 4 for  $\alpha = 0.4$  as representative figures. Similar dynamics is observed for  $\alpha = 0.6$  and  $\alpha = 0.8$ .

![](_page_5_Figure_1.jpeg)

**Figure 4.** We plot  $x(t-1)$  versus  $x(t)$  for Lozi map after discarding several transients for fractional order ( $\alpha = 0.4$ ) Lozi map for few values of  $\beta$ 

### **3. Results and discussion**

From the bifurcation diagrams, it is clear that the fractional-order parameter  $\alpha$  affects the evolution of the fractional Lozi map system and the system can become unstable in the range the original map was stable. However, the overall bifurcation structure does not change much even for small values of  $\alpha$ . The map parameter  $\beta$  also affects the evolution of the system. An increase in  $\beta$  however makes the system more unstable for the same value of  $\alpha$ . We can explore larger values of β for larger values of  $\alpha$ . From the control viewpoint, we find that the stable fixed point is observed over a larger range of parameters for smaller  $\alpha$ . However, there is no significant change in the parameter regime where chaos is observed as we change  $\alpha$ . Thus in this particular case, the fractional version of the map can be very useful for stabilizing the fixed point. The robustness of chaos is not lost and there are no periodic orbits inside the chaotic domain. In the

chaotic region, the fractional order map is effectively higher dimensional compared to its integer order version. Hence it can have applications in cryptography and secure communication [30].

There have been several analytical results in the Lozi map. The genuine Lozi system does not show a period-doubling cascade. The smoother versions of this map show a period-doubling cascade. They also show abrupt transitions. We may think that the fractional version of the map is smoother and qualitatively different behaviour is obtained. This is not the case. Our results indicate that the dynamics deforms continu[ousl](#page-7-8)y on the introduction of the memory. It would be interesting to study if similar analytic studies can be carried out for the fractional Lozi map.

#### **Acknowledgement**

PMG thanks DST-SERB (CRG/2020/003993) for financial assistance.

# **Conflict of interest**

The authors declare no competing financial interest.

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