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# **Applications of the Nonstandard Finite Difference Method to a Fractional Model Explaining Diabetes Mellitus and Its Complications**

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**Abstract:** This work examines a mathematical model of high blood sugar conditions and its consequences in a society using fractional differential equations. It attempts to solve the problem using a nonstandard way because standard finite difference numerical methods can result in numerical instabilities. The nonstandard finite difference scheme (NSFDS), which satisfies dynamical consistency, is the recommended nonstandard method for discretising the model. To show that the model is stable at the equilibrium positions, analyses of both discrete and continuous models are studied.Stability analysis is carried out at the discretised model's equilibrium point using the Schur-Cohn criterion. Consequently, the model's asymptotically stable state is demonstrated. Furthermore, by contrasting the stability for various step sizes with conventional techniques like Finite Difference Scheme (FDS), the benefits of the NSFDS are shown. The NSFDS has been shown to converge at bigger step sizes. Furthermore, a graphical comparison is shown between the numerical findings acquired by the NSFDS and the FDS. It is noted that the NSFDS is accurate.

*Keywords***:** diabetes mellitus, nonstandard finite difference scheme, stability analysis, fractional calculus

**MSC:** 34A30, 65L07

# **Abbreviation**

NSFDS Nonstandard finite difference scheme FDS Finite difference scheme

# **1. Introduction**

Differential equations can be applied to model a diversity of biological issues. As is well known, diabetes has recently become a highly prevalent condition. Numerous significant studies on diabetes have been conducted. Stability analysis

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plays a crucial part in epidemic models. The following is a summary of some diabetes research: In order to provide a higher quality for people, the authors in [1] suggest a mathematical model of high blood sugar. They provided the linear model's numerical solution and stability analysis, where the unknowns are the proportion of diabetics with and without issues. In their research, Akinsola and Oluyo [2–4] used various techniques to arrive at the numerical and analytical solution of the diabetes mellitus complications and control model. Furthermore, AlShurbaji et al. [5] consider the linear diabetes mellitus model. Numerical techniques inc[lu](#page-20-0)ding Euler, Heun, Runge-Kutta, and Adams-Moulton were applied to get the solution of a system of linear differential equations. Additionally, the Euler-Cauchy approach was used to offer a numerical solution by Vanitha and Porchelvi [6]. By usi[ng](#page-20-1) [th](#page-20-2)e right parameters, the authors in [7] constructed a nonlinear mathematical model of high blood sugar conditions. Numerical experiments and stability analysis are [pr](#page-20-3)esented. Global asymptotic stability and the model suggested in [7] were taken into consideration by De Oliveira et al. [8]. Numerical simulations were used to confirm stability. The dynamics of a population of healthy individuals, diabetics and pre-diabetics with and without problems were given by a[ut](#page-20-4)hors in [9]. The theory of best favorable is appl[ie](#page-20-5)d. Controllability and stability for both linear and nonlinear differential equations were demonstrated by Permatasari et al. [10], who worked on the model created in [9]. The nonlinear diabetic [me](#page-20-5)llitus model was examined by Widyaningsih et al. [[11](#page-21-0)] in relation to genetic and lifestyle factors. The number of deaths from diabetes was predicted using the 4th-order Runge-Kutta method. A model that is linear which describes diabetes me[lli](#page-21-1)tus and its consequences was the subject of stability investigated by Aye [12]. The Bellman and Coke theorem was used to test the stability.Aye et al. [13] used th[e H](#page-21-2)omotopy Perturbation Method to solve [a c](#page-21-1)omparable scenario. Aye [14] investigated how to control affected the same mo[del](#page-21-3).

It is well known that the fields of applied mathematics heavily rely on stability analysis of mathematical models. In stability analysis, discretizing the models is crucial in real-world scenarios. Numerical instabilities might a[lso](#page-21-4) result from using typical numerical methods to solve such situations. N[ons](#page-21-5)tandard approaches are therefore important. The nonstandard finite differenc[e sc](#page-21-6)heme (NSFDS), created by Mickens [15–20], is one of the most practical and efficient discrete approaches. Additionally, it offers convergence outcomes at even larger step sizes. The studies of Patidar [21, 22] give a thorough literature review on NSFDS. Numerous applied mathematics disciplines have produced numerous studies on NSFDS. The following is a list of some recent research on NSFDS and stability analysis: A mathematical model based on the Tacoma Narrows Bridge collapse is examined by Adek[any](#page-21-7)[e a](#page-21-8)nd Washington [23]. For the torsional and vertical variants, two NSFDS are built. Anguelov et al. [24] proposed the applicability of NSFDS to Ebola virus [sys](#page-21-9)[tem](#page-21-10) of model in Africa. Arenas et al. [25] introduced fractional model of epidemic for susceptible-infected-recovered (SIR) and susceptible infected (SI). There are some similarities with classical approaches after applying NSFDS. For integer and fractional order, a unique chaotic model was analyzed by Baleanu et al. [26]. In both [case](#page-21-11)s stability analysis was presented. NSFDS is utilized to present numerical simula[tion](#page-21-12)s. NSFDS are constructed for two metapopulation models by Dang and Hoang [27]. Other featu[res](#page-21-13) including positivity, boundedness, and monotone convergence were provided along with stability analysis. The theoretical investigation was supported by numerical computations. Kocabıyık et al. [28] and Dang and Hoang [29] used the NSFDS to approximate a computer virus model[. U](#page-21-14)sing NSFDS, the authors in [30] build a mathematical system that describes the Michaelis-Menten harvesting rate. Kocabiyik and Ongun [31] described the NSFDS of smoki[ng m](#page-21-15)odel in discretization way to ascertain the impacts on humans who smoke. In the study by Ongun and Arslan [32], two distinct NSFDS for the fractional order Hantavirus model were compared. Shabbir et al. [[33\]](#page-21-16) used NSFDS to build a [pre](#page-21-17)dator-prey model. Investigations are conducted into stability analysis as well as other char[acte](#page-21-18)ristics including positivity, boundedness, and persistence of solutions. An NSFDS for the SICA-AIDS mode[l w](#page-22-0)as put forth by Vaz and Torres [34]. Both global and elementary stability were examined. In [35], Egbelowo et al. examined a linear mathematic[al m](#page-22-1)odel of pharmacokinetics. The analytical solution, NSFDS, and Finite Difference Scheme (F[DS\)](#page-22-2) were shown. [36–40, 40, 41] give more recent research on the stability study of the mathematical models. Said et al. [42] presented a NSFDS for solving a fractional decay model. The proposed NSFDM was constructed by incorporating a non-standard de[nom](#page-22-3)inator function, resulting in an explicit numerical scheme a[s ea](#page-22-4)sy as the conventional Euler method, but it provides very accurate solutions and has unconditional stability. It was found that the method's estimated errors are extremel[y m](#page-22-5)[inimal,](#page-22-6)s[uch](#page-22-7) as within the machine precision. Iikem Cetinkaya [43] presented and described the complicati[ons](#page-22-8) of diabetes mellitus and its consequences, concluded that the model is asymptotically stable.

Extension of ordinary calculus to fractional calculus takes us into a new mathematical domain for solution of diverse problems in various fields of study, such as computational biology, probability, control theory etc. Differential equations can represent these phenomena in a better way. The properties of well-known fractional calculus were used to get solutions of certain fractional calculus equations representing these types of phenomena. The solution of most of these differ-integral equations take the form of Mittag-Leffler function, Fox-Wright functions etc. There are numerous authors who define fractional calculus. The Riemann-Liouvellie (R-L) definition is the one that is most frequently employed [44–48]. The Caputo definition of fractional derivative (1967) [44–48] is another helpful formulation. To prevent nonzero fractional derivative of a constant function, applying left hand Jumarie's version of the R-L fractional derivative is helpful [49, 50]. A theory of characterizing the points where the function is non-differentiable, using a Jumarie type right hand fractional derivative version of the R-L was recently developed by Ghosh et al. in the publication [51]. Different type[s of](#page-22-9) [sol](#page-22-10)utions were provided by differential equations in various [fra](#page-22-9)[ctio](#page-22-10)nal derivative forms [44–48]. Therefore, fractional differential equations cannot be solved using a normal algorithm. Therefore, a growing area of applied mathematics in fr[acti](#page-22-11)[ona](#page-22-12)l calculus gives the solution and interpretation of such models. Recently, the Predictor-Corrector approach [52], Adomain decomposition method [45, 53, 54], Disordered Perturbation method [55], Iteration metho[d th](#page-22-13)rough variation [56], Method through differential transform [57] were employed to get the solution of both li[nea](#page-22-9)[r an](#page-22-10)d non-linear differential equations. An analytical approach for solving linear FDE's with the derivatives by Jumarie [50] using sine, cosine and Mittag-Leffler functions was recently devised by Ghosh et al. [58]. The numerical analysis of the nuclear decay model [and](#page-22-14) fractional population growth was [pres](#page-22-15)[ent](#page-22-16)[ed b](#page-22-17)y Alzaid et al. [59]. By applying e[xpo](#page-22-18)nential decay kernel, Alqhtani et al. [[60\]](#page-22-19) provided solutions of space-fractional di[ffus](#page-22-20)ion equations numerically.

NSFDS have a particular advantage over traditional finite difference sc[hem](#page-22-12)es because they may provide exact numerical schemes that lead to exact solutions [at t](#page-23-0)he nodal points. For the purpose of solving certain classes of firstorder differential equations, such as the logistic [an](#page-23-1)d exponential growth, systems of linear ODE models, [hig](#page-23-2)her order ODEs, some PDE problems, etc., Micken was the first to design and explain the derivation of exact finite difference techniques [61, 62]. Initially, NSFDS were designed to offer either an accurate finite difference scheme or the optimal numerical scheme, with the goal of achieving all the solutions qualitatively of the original problem.

The purpose of this piece of work is to design a numerical scheme employing nonstandard finite difference discretization to solve the diabetes model and its consequences given in linear fractional differential equations. It compares the perform[anc](#page-23-3)[e o](#page-23-4)f the proposed scheme to other methods found in the literature. The main features of the proposed numerical approach are its explicit Euler method-like design, high order of convergence, and unconditional stability.

The paper is organized in the following manner:

Section 1 provides an insight of fractional calculus and the development of an NSFDS. The mathematical preliminaries are given in section 2. In the Section 3 proposed methodology is introduced. Section 4 employs numerical examples from existing literature to evaluate the validity of the proposed technique. Finally, the discussion and conclusions are presented in Section 5.

### **2. Mathematical preliminaries**

#### **Definition 2.1 Riemann-Liouville (R-L) fractional derivative:**

The fractional Riemann-Liouville (R-L) left derivative is given by,

$$
{}_{b}D_{z}^{\alpha}f(z) = \frac{1}{\Gamma(n+1-\alpha)} \left(\frac{d}{dz}\right)^{n+1} \int_{b}^{z} (z-\tau)^{n-\alpha} f(\tau) d\tau,\tag{1}
$$

where  $n \le \alpha < n+1$ , *n* is positive integer.

When  $0 \leq \alpha < 1$  then

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$$
bD_z^{\alpha} f(z) = \frac{1}{\Gamma(1-\alpha)} \frac{d}{dz} \int_b^z (z-\tau)^{-\alpha} f(\tau) d\tau.
$$
 (2)

The corresponding fractional Riemann-Liouville (R-L) right derivative is given as follows:

$$
{}_{z}D_{b}^{\alpha}f(z) = \frac{1}{\Gamma(1-\alpha)} \frac{-d}{dz} \int_{z}^{b} (\tau - z)^{m-\alpha} f(\tau) d\tau,
$$
\n(3)

where  $n \leq \alpha \leq n+1$ .

According to the definitions (1)-(3) above, a constant's derivative is non-zero, which is in opposition to the constant's classical derivative of zero. A revision to the R-L concept of fractional derivative was offered by Caputo [29] in 1967 in order to get around this drawback.

#### **Definition 2.2 Caputo Derivative:**

Fractional derivative defined by Caputo is given in the following form [49]:

$$
{}_{b}^{C}D_{z}^{\alpha}f(z) = \frac{1}{\Gamma(m-\alpha)} \int_{b}^{z} (z-\tau)^{m-\alpha-1} f^{m}(\tau) d\tau,
$$
\n(4)

where  $m-1 \leq \alpha < m$ .

In this concept,  $f(x)$  is differentiated *n* times first, and then  $n - \alpha$  times are integrated. This method's drawback is that *f*(*x*) needs to be *n* times differentiable for the existence of α-th order, where  $m − 1 ≤ α < m$ . This definition is not applicable if the function is non-differentiable. This method has two key advantages:

(i) the derivative of a constant is zero in fractional way; and

(ii) the initial conditions of the R-L type differential equations are fractional type, whereas the Caputo type fractional differential equations have classical derivative type initial conditions.

This indicates that a fractional differential equation including R-L fractional derivatives necessitates an understanding of fractional initial conditions, which can occasionally be challenging to physically interpret [45].

#### **Definition 2.3 Jumaire type fractional derivative:**

In order to get around the fractional derivative of a constant that is not zero, Jumarie [50] suggested another way to define the left R-L type fractional derivative in the interval  $[a, b]$  of the function  $f(x)$  in the f[orm](#page-22-15) as follows:

$$
f_L^{(\alpha)}(z) =_d^l D_z^{\alpha} f(z) \tag{5}
$$

$$
= \begin{cases} \frac{1}{\Gamma(-\alpha)} \int_a^z (z-\tau)^{-\alpha-1} f(\tau) d\tau, & \alpha < 0 \\ \\ \frac{1}{\Gamma(1-\alpha)} \frac{d}{dz} \int_a^z (z-\tau)^{-\alpha} (f(\tau) - f(a)) d\tau, & 0 < \alpha < 1 \\ \\ \left( f^{(\alpha-n)}(z) \right)^{(n)}, & n \le \alpha < n+1. \end{cases}
$$

For  $z < a$ , they assume that  $f(z) - f(a) = 0$ . The R-L derivative of order  $\alpha$  ( $0 < \alpha < 1$ ) of the offset function  $(f(x) - f(a))$  is shown in the second line of equation (5), while the first expression represents fractional integration. They employ the third line for  $\alpha > 1$ , which means that they apply complete n order differentiation to the offset function with order  $0 < (\alpha - n) < 1$  after first differentiating it using the formula of the second line. In this case, they selected the integer n, which is just less than the real number  $\alpha$  ( $n \leq \alpha < n+1$ ).

#### **Definition 2.4 Mittag-Leffler function:**

This function was defined and explored by Mittag-Leffler. The exponential function,  $e^x$ , is directly generalized by this function, which is an essential part of fractional calculus. Power series expansion has been used to define the one, two, and three-parameter Mittag-Leffler functions as follows [63–65].

$$
E_{\alpha}(x) = \sum_{i=0}^{\infty} \frac{x^{i}}{\Gamma(1+\alpha i)}, \text{ for } \alpha \in C, Re(\alpha) > 0.
$$
 (6)

$$
E_{\alpha, \beta}(x) = \sum_{i=0}^{\infty} \frac{x^i}{\Gamma(\beta + \alpha i)}, \text{ for } \alpha, \beta \in C, Re(\alpha) > 0.
$$
 (7)

$$
E_{\alpha, \beta}^{\gamma}(x) = \sum_{i=0}^{\infty} \frac{(\gamma)_{i}x^{i}}{\Gamma(\beta + \alpha i)i!}, \text{ for } \alpha, \beta, \gamma \in C, Re(\alpha) > 0.
$$
 (8)

where,  $(\gamma)_i = \gamma(\gamma + 1)(\gamma + 2)(\gamma + 3)...(\gamma + i - 1)$  and  $(\gamma)_0 = 1$ .

The representation in terms of integral is given in [63, 66–68] as follows,

$$
E_{\alpha}(z)=\frac{1}{2\pi}\int_{C}\frac{t^{\alpha-1}e^{t}}{t^{\alpha}-z}dt, \ z\in\mathbb{C}, \ Re(\alpha)>0.
$$

In this case, the path of integration *C* is a loop that encloses the disk's circles  $|t| \leq |z|^{1/\alpha}$  in the positive sense: *|arg*(*t*)*| ≤* <sup>π</sup> on *C* [63, 66–68]. It begins and ends at *−*∞.

The two parameter Mittag-Leffler function's appropriate integral representation is provided by [63, 66–68] as follows:

$$
E_{\alpha, \beta}(z) = \frac{1}{2\pi} \int_C \frac{t^{\alpha-\beta} e^t}{t^{\alpha}-z} dt, \ z \in \mathbb{C}, \ Re(\alpha) \text{ is } +ve.
$$

Where *C*, the contour remains same as already defined above.

# **3. Modelling of the fractional diabetes mellitus system**

The non-fractional model is given in [43] as below:

$$
\begin{cases}\n\frac{dH}{dt} = -(\tau + \mu)H + \beta \theta + \sigma S, \\
\frac{dS}{dt} = -(\mu + \alpha_p + \beta(1 - \theta) + \sigma)S + \tau H, \\
\frac{dD}{dt} = \omega T + \alpha_p S - (\lambda_p + \mu)D, \\
\frac{dC}{dt} = -(\mu + \delta + \gamma)C + \lambda_p D, \\
\frac{dT}{dt} = -(\mu + \omega)T + \gamma C,\n\end{cases}
$$

such that the initial values at time  $t = 0$  are given by  $H_0$ ,  $S_0$ ,  $D_0$ ,  $C_0$ , and  $T_0$ . The state variables  $H(t)$ ,  $S(t)$ ,  $D(t)$ ,  $C(t)$ , and *T*(*t*) represent totally healthy, suspected, confirmed without complications, confirmed with complications and confirmed with complications receiving treatment, respectively. The parameters γ, τ, θ, α<sub>*p*</sub>, σ, ω, λ<sub>*p*</sub>, μ, δ, and β, denote birth rate of children, healthy born children rate, naturally caused mortality rate, healthy population becoming suspected population rate, suspected population becoming healthy population rate, diabetes incidence rate probability, confirmed diabetic population developing complications rate, confirmed diabetic with complexity becoming without complexity rate, natural mortality with complexities rate and cured population rate from diabetes.

As an extension of the system above, the following diabetes mellitus mathematical system is considered using the fractional derivatives of order <sup>α</sup>. A system of linear differential equations with fractional order <sup>α</sup> is considered as:

$$
\begin{cases}\n^c D_t^\alpha H = -(\tau + \mu)H + \beta \theta + \sigma S, \\
^c D_t^\alpha S = -(\mu + \alpha_p + \beta(1 - \theta) + \sigma)S + \tau H, \\
^c D_t^\alpha D = \omega T + \alpha_p S - (\lambda_p + \mu)D, \\
^c D_t^\alpha C = -(\mu + \delta + \gamma)C + \lambda_p D, \\
^c D_t^\alpha T = -(\mu + \omega)T + \gamma C,\n\end{cases} \tag{9}
$$

such that the initial values are same as stated above in the non-fractional model.

An autonomous vector field of a general form can be considered as below:

$$
Y'=S(Y), Y \in \mathbb{R}^n,
$$
\n<sup>(10)</sup>

the corresponding fractional system after linearization can be defined as

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 ${}^{c}D_{t}^{\alpha}\psi=J(E^*)\psi,$ 

such that the  $E^*$  represents the equilibrium points of the corresponding steady system given in 10, and  $J(E^*)$  represents Jacobian of the same system above. Negative real values are assumed from the Jacobian establishing the asymptotic stability of the system 10, as in [40].

# **4. Diabetes mellitus f[rac](#page-22-6)tional system discretization by NSFDS**

The fractional order diabetes mellitus system is discretized using the non-standard differences. On the basis of the guidelines in [16], the construction of the numerical scheme is carried out using non-standard differences. The development of the weight functions relies on the eigen-values of the matrix obtained from the fractional system.

• To counter the instabilities arising numerically, the discretized differential should be same as the differentials in the given system.

• A general fo[rm](#page-21-19) representing fractional discretized derivatives of order  $\alpha$  is given below:

$$
{}^cD_t^{\alpha}x \to \frac{x_{n+1} - \psi(h)x_n}{\phi(h)},
$$

where  $\psi(h, \alpha, p)$  and  $\phi(h, \alpha, p)$ , where *p* is the set of model's parameters, are called the weight functions.

• Nonlocal terms are constructed in a discrete setting to replace the nonlinearities such that  $s^2 \to s_{n-1}^2$  and  $s^2 \to s_n s_{n-1}$ .

• Furthermore, the corresponding difference model should be aligned with the original differential model in terms of the conditions imposed.

The positivity requirement from above is satisfied as the system is discretized in the following setup:

(i) For  $H(t)$ :  $H(t) \rightarrow H(n+1)$  and  $S(t) \rightarrow S(n)$  are used.

(ii) For  $S(t)$  and  $H(t)$ :  $H(t) \rightarrow H(n)$  and  $S(t) \rightarrow S(n+1)$  are used.

(iii) For  $S(t)$ ,  $D(t)$ , and  $T(t)$ :  $S(t) \rightarrow S(n)$ ,  $D(t) \rightarrow D(n+1)$ , and  $T(t) \rightarrow T(n)$  are used.

(iv) For  $D(t)$  and  $C(t)$ :  $D(t) \rightarrow D(n)$  and  $C(t) \rightarrow C(n+1)$  are used.

(v) For  $C(t)$  and  $T(t)$ :  $C(t) \rightarrow C(n)$  and  $T(t) \rightarrow T(n+1)$  are constructed.

Thus, the following discrete model is obtained:

$$
\begin{cases}\nH(n+1) = \frac{H(n) + (\beta \theta + \sigma S(n))\phi_1}{1 + (\mu + \tau)\phi_1}, \\
S(n+1) = \frac{S(n) + (\beta(1-\theta) + \tau H(n))\phi_2}{1 + (\mu + \alpha_p + \sigma)\phi_2}, \\
D(n+1) = \frac{D(n) + (\alpha_p S(n) + \omega T(n))\phi_3}{1 + (\mu + \lambda_p)\phi_3}, \\
C(n+1) = \frac{C(n) + \lambda_p \phi_4 D(n)}{1 + (\mu + \delta + \gamma)\phi_4}, \\
T(n+1) = \frac{T(n) + \gamma \phi_5 C(n)}{1 + (\mu + \omega)\phi_5},\n\end{cases} (11)
$$

such that  $\phi_i$ ,  $i = 1 \rightarrow 5$  represents the corresponding weight functions to all the state variables given below as:

$$
\begin{cases}\n\phi_1 = \frac{-E_{\alpha}(-(\mu + \tau)h) + 1}{(\mu + \tau)}, \\
\phi_2 = \frac{-E_{\alpha}(-(\mu + \alpha_p + \sigma)h) + 1}{(\mu + \alpha_p + \sigma)}, \\
\phi_3 = \frac{-E_{\alpha}(-(\mu + \lambda_p)h) + 1}{(\mu + \lambda_p)}, \\
\phi_4 = \frac{-E_{\alpha}(-(\mu + \delta + \gamma)h) + 1}{(\mu + \delta + \gamma)}, \\
\phi_5 = \frac{-E_{\alpha}(-(\mu + \omega)h) + 1}{(\omega + \mu)}.\n\end{cases}
$$

The standard finite differences carry the approximation errors over the iterations resulting in instabilities as numerical computations are performed. The above discretization using non-standard differences helps achieve better approximations as the construction and the consequent presence of the weight functions based on eigen-value decomposition incorporates better corrections as compared to the standard classical schemes such as Euler and Runge-Kutta discretizations.

#### **4.1** *Analysis*

From the above model (9) we write in the matrix-vector formulation as follows:

$$
(X^{\alpha})^{'} = AX + b,\tag{12}
$$

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where,  $X^{\alpha} = [H^{\alpha}, S^{\alpha}, D^{\alpha}, C^{\alpha}, T^{\alpha}]^T$ ,

$$
A = \left[\begin{array}{cccccc}-(\tau+\mu) & \sigma & 0 & 0 & 0 \\-(-\tau) & -(\alpha+\sigma+\mu) & 0 & 0 & 0 \\ 0 & \alpha & -(\lambda+\mu) & 0 & \omega \\ 0 & 0 & \lambda & -(\gamma+\delta+\mu) & 0 \\ 0 & 0 & 0 & \gamma & -(\omega+\mu)\end{array}\right],
$$

and

$$
b = \left[\begin{array}{c} \beta\theta \\ \beta(1-\theta) \\ 0 \\ 0 \\ 0 \end{array}\right].
$$

We decompose the matrix *A* such that equation (12) becomes

$$
\left(X^{\alpha}\right)' = P^{-1}DP + b.\tag{13}
$$

Where *D* is the matrix of the eigen values and *P* is the matrix of the eigen vectors such that

 $P = [w_1, w_2, w_3, w_4, w_5].$ 

We write the solution using the Mittag-Leffler function such that

$$
U(t) = P \begin{bmatrix} c_1 E_{\alpha}(\lambda_1 t^{\alpha}) \\ c_2 E_{\alpha}(\lambda_2 t^{\alpha}) \\ c_3 E_{\alpha}(\lambda_3 t^{\alpha}) \\ c_4 E_{\alpha}(\lambda_4 t^{\alpha}) \\ c_5 E_{\alpha}(\lambda_5 t^{\alpha}) \end{bmatrix} + \begin{bmatrix} \beta \theta \\ \beta(1 - \theta) \\ 0 \\ 0 \\ 0 \end{bmatrix},
$$

or

$$
U(t) = [w_1, w_2, w_3, w_4, w_5] \begin{bmatrix} c_1 E_{\alpha}(\lambda_1 t^{\alpha}) \\ c_2 E_{\alpha}(\lambda_2 t^{\alpha}) \\ c_3 E_{\alpha}(\lambda_3 t^{\alpha}) \\ c_4 E_{\alpha}(\lambda_4 t^{\alpha}) \\ c_5 E_{\alpha}(\lambda_5 t^{\alpha}) \end{bmatrix} + \begin{bmatrix} \beta \theta \\ \beta(1 - \theta) \\ 0 \\ 0 \\ 0 \end{bmatrix},
$$

where,

$$
U(0) = [w_1, w_2, w_3, w_4, w_5] \begin{bmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \\ c_5 \end{bmatrix} + \begin{bmatrix} \beta \theta \\ \beta(-\theta + 1) \\ 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} u_1(0) \\ u_2(0) \\ u_3(0) \\ u_4(0) \\ u_5(0) \end{bmatrix},
$$

and

$$
E_{\alpha}(\lambda t^{\alpha}) = \sum_{n=0}^{\infty} \frac{\lambda^{\alpha}(t)^{\alpha n}}{\Gamma(\alpha n + 1)}.
$$

# **5. Discretized fractional diabetes mellitus system stability**

The stability of the discretized model is discussed in this section. The Lemma below outlines a special case when all the state variables, after incorporating the initial states, are all non-negative. We present the Lemma as:

**Lemma 1** If the following inequality:

$$
-((-\theta+1)\beta+\tau H(n))\left\langle \frac{S(n)}{\phi_2}\right\rangle,
$$

holds for *S*, *H*,  $\beta$ ,  $\theta$ ,  $\tau$  and  $\phi$ . Then all analytical representations of the state variables in (11) are positive along with the positive initial states.

The Lemma 1 has been proved in [43] by establishing the above inequality and the positive weight functions.

A local asymptotic stability result is presented below.

**Theorem 5.1** [43] Suppose *J* is the Jacobian and *E ∗* represents equilibrium points of the fractional system. Then the fractional system demonstrates asymptotic stability if real parts from  $J(E^*)$  are all negative.

**Proof.** The equilibrium point of s[yste](#page-22-21)m 9 is obtained as  $E^* = (H^*, S^*, D^*, C^*, T^*)$ , where

$$
H^* = \frac{\beta \theta}{\mu + \tau} + \left( \frac{\beta(1-\theta) + \frac{\beta \theta \tau}{\mu + \tau}}{(\mu + \alpha + \sigma) - \frac{\tau \sigma}{\mu + \tau}} \right) \frac{\sigma}{\mu + \tau},
$$

$$
S^* = \frac{\beta(1-\theta) + \frac{\beta \theta \tau}{\mu + \tau}}{(\mu + \alpha + \sigma) - \frac{\tau \sigma}{\mu + \tau}},
$$

$$
D^* = \frac{\alpha\beta(1-\theta) + \alpha\beta\theta\tau}{(\mu+\lambda)[(\mu+\alpha+\sigma)(\mu+\tau)-\tau\sigma]} + \frac{\left[\frac{\gamma\lambda\alpha\beta(1-\theta) - \frac{\gamma\lambda\alpha\beta\theta\tau}{\mu+\tau}}{-(\mu+\delta+\lambda)(\mu+\lambda)[(\mu+\alpha+\sigma) - \frac{\tau\sigma}{\mu+\tau}]}\right]}{-(\mu+\omega) + \frac{\gamma\lambda\omega}{(\mu+\lambda)(\mu+\delta+\gamma)}} \times \left[\frac{\omega}{\mu+\lambda}\right],
$$

$$
C^* = \frac{-\lambda \alpha \beta (1-\theta) + \frac{\lambda \alpha \beta \theta \tau}{\mu + \tau}}{-(\mu + \delta + \lambda)(\mu + \lambda)[(\mu + \alpha + \sigma) - \frac{\tau \sigma}{\mu + \tau}]}
$$

$$
+\frac{\left[\dfrac{\gamma\lambda\alpha\beta(1-\theta)-\dfrac{\gamma\lambda\alpha\beta\theta\tau}{\mu+\tau}}{-(\mu+\delta+\lambda)(\mu+\lambda)[(\mu+\alpha+\sigma)-\dfrac{\tau\sigma}{\mu+\tau}]}\right]}{-(\mu+\omega)+\dfrac{\gamma\lambda\omega}{(\mu+\lambda)(\mu+\delta+\gamma)}}\times\Bigg[\dfrac{\lambda\omega}{(\mu+\lambda)(\mu+\delta+\gamma)}\Bigg],
$$

$$
T^* = \frac{\left[\frac{\gamma\lambda\alpha\beta(1-\theta) - \frac{\gamma\lambda\alpha\beta\theta\tau}{\mu+\tau}}{-(\mu+\delta+\lambda)(\mu+\lambda)\left[(\mu+\alpha+\sigma) - \frac{\tau\sigma}{\mu+\tau}\right]}\right]}{-(\mu+\omega) + \frac{\gamma\lambda\omega}{(\mu+\lambda)(\mu+\delta+\gamma)}}.
$$

The Jacobian matrix of the continuous model at the equilibrium point  $E^* = (H^*, S^*, D^*, C^*, T^*)$  determined as

$$
J(H^*,\ S^*,\ D^*,\ C^*,\ T^*) = \left[ \begin{array}{cccccc} -(\mu+\tau) & \sigma & 0 & 0 & 0 \\ \tau & -(\mu+\alpha+\sigma) & 0 & 0 & 0 \\ 0 & \alpha & -(\mu+\lambda) & 0 & \omega \\ 0 & 0 & \lambda & -(\mu+\delta+\gamma) & 0 \\ 0 & 0 & 0 & \gamma & -(\mu+\omega) \end{array} \right].
$$

Locally asymptotic stability of the fractional model can be analyzed by obtaining the eigenvalues of the Jacobian matrix at equilibrium points. Since all the eigen values of the Jacobian matrix are negative real numbers then the system is asymptotically stable.

We introduce the Schur-Cohn criterion for our system.

**Lemma 2** [41] Consider the characteristic polynomial

$$
q(\lambda) = \lambda^5 + a_1 \lambda^4 + a_2 \lambda^3 + a_3 \lambda^2 + a_4 \lambda + a_5,
$$

where  $b_1$ ,  $b_2$ ,  $b_3$ ,  $b_4$ ,  $b_5$  are constants. The zeros of the characteristic polynomial defined by above lie inside the unit disk if and only if the following conditions hold:

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a.  $q(1) > 0$ , **b.** 1−*b*<sub>1</sub> +*b*<sub>2</sub> −*b*<sub>3</sub> +*b*<sub>4</sub> −*b*<sub>5</sub> > 0, c. The matrices

$$
Q_4^{\pm} = \left[\begin{array}{cccc} 1 & 0 & 0 & 0 & 0 \\ b_1 & 1 & 0 & 0 & 0 \\ b_2 & b_1 & 1 & 0 & 0 \\ b_3 & b_2 & b_1 & 1 & 0 \\ b_4 & b_3 & b_2 & b_1 & 1 \end{array}\right] \pm \left[\begin{array}{cccc} 0 & 0 & 0 & 0 & b_5 \\ 0 & 1 & 0 & b_5 & b_4 \\ 0 & 0 & b_5 & b_4 & b_3 \\ 0 & b_5 & b_4 & b_3 & b_2 \\ b_5 & b_4 & b_3 & b_2 & b_1 \end{array}\right],
$$

are positive.

**Proof.** One can conclude that if the Schur-Cohn criterion is satisfied, then the system is asymptotically stable. Therefore, the Jacobian matrix at the equilibrium point  $E^* = (H^*, S^*, D^*, C^*, T^*)$  can be written as:

$$
J^* = J(H^*, S^*, D^*, C^*, T^*) = \begin{bmatrix} h_1 & \phi_1 \sigma h_1 & 0 & 0 & 0 \\ \phi_2 \tau h_2 & h_2 & 0 & a_5 & b_4 \\ 0 & \phi_3 \alpha h_3 & h_3 & 0 & \phi_3 \omega h_3 \\ 0 & 0 & \phi_4 \lambda h_4 & h_4 & 0 \\ 0 & 0 & 0 & \phi_5 \gamma h_5 & h_5 \end{bmatrix},
$$

where,

$$
h_1 = \frac{1}{1 + \phi_1(\mu + \tau)},
$$
  
\n
$$
h_2 = \frac{1}{1 + \phi_2(\mu + \alpha + \sigma)},
$$
  
\n
$$
h_3 = \frac{1}{1 + \phi_3(\mu + \lambda)},
$$
  
\n
$$
h_4 = \frac{1}{1 + \phi_4(\mu + \delta + \gamma)},
$$
  
\n
$$
h_5 = \frac{1}{1 + \phi_5(\mu + \omega)},
$$

where the coefficients are determined as

$$
b_1 = -(h_1 + h_2 + h_3 + h_4 + h_5),
$$
  
\n
$$
b_2 = (h_3 + h_4)h_5 + h_3h_4 + (h_1 + h_2)(h_3 + h_4 + h_5) - (1 - \tau \sigma \phi_1 \phi_2)h_1 h_2,
$$
  
\n
$$
b_3 = -(1 + \gamma \lambda \omega \phi_3 \phi_4 \phi_5)h_3 h_4 h_5 + (\tau \sigma \phi_1 \phi_2 - 1)h_1 h_2 (h_3 + h_4 + h_5) - (h_1 + h_2)[h_5 (h_3 + h_4) + h_3 h_4],
$$
  
\n
$$
b_4 = (h_1 + h_2)(1 + \phi_3 \phi_4 \phi_5 \lambda \omega \gamma) + h_1 h_2 (1 - \sigma \tau \phi_1 \phi_2) [h_5 (h_3 + h_4) + h_3 h_4],
$$
  
\n
$$
b_5 = h_1 h_2 h_3 h_4 h_5 (\sigma \tau \phi_1 \phi_2 - 1)(1 + \phi_3 \phi_4 \phi_5 \lambda \omega \gamma),
$$

we find the coefficients of the characteristic polynomial as below:

*b*<sup>1</sup> = *−*4*.*95460854*, b*<sup>2</sup> = 9*.*97630952*, b*<sup>3</sup> = *−*9*.*9657420912*, b*<sup>4</sup> = 4*.*9873210654*, b*<sup>5</sup> = *−*0*.*98322055324*.*

In the view of Lemma 2, since

a.  $q(1) = 1 + b_1 + b_2 + b_3 + b_4 + b_5 = 0.213 \times 10^{-12} > 0$ , b. *−p*(*−*1) = 1*−b*<sup>1</sup> +*b*<sup>2</sup> *−b*<sup>3</sup> +*b*<sup>4</sup> *−b*<sup>5</sup> = 30*.*78648021 *>* 0,

c. The matrices

$$
Q_4^{\pm} = \left[ \begin{array}{cccc} 1 & 0 & 0 & \pm b_5 \\ b_1 & 1 & \pm b_5 & \pm b_4 \\ b_2 & b_1 \pm b_5 & 1 \pm b_4 & \pm b_3 \\ b_3 \pm b_5 & b_2 \pm b_4 & b_1 \pm b_3 & 1 \pm b_2 \end{array} \right].
$$

As all the axioms of Schur-Cohn are established, the given system is asymptotically stable locally for the chosen parameters.

# **6. Numerical experiment**

This section represents numerical experiment in order to make analytical results, standard finite differences and the non-standard finite difference approximations. The parameter values are chosen such that  $θ = 0.913$ ,  $β = 0.037$ ,  $τ =$ 0.035,  $\mu = 0.51$ ,  $\sigma = 0.081$ ,  $\lambda_p = 0.512$ ,  $\alpha_p = 0.0231$ ,  $\delta = 0.02001$ ,  $\omega = 0.079$ , and  $\gamma_p = 0.081$ . The eigen values are calculated using the matrix *A* given in (12).

$$
\begin{pmatrix}\n\lambda_1 \\
\lambda_2 \\
\lambda_3 \\
\lambda_4 \\
\lambda_5\n\end{pmatrix} = \begin{pmatrix}\n-0.596819223081673 \\
-0.310501484379986 \\
-0.252031242374329 \\
-0.123968757625671 \\
-0.126679292538342\n\end{pmatrix}.
$$

Since all the eigen values are negative real numbers, Theorem 5*.*1 dictates the asymptotic stability of the fractional system locally. The corresponding eigen vectors are given as follows:

$$
w_1 = \begin{pmatrix} 5.896221210428095 \times 10^{-17} \\ 1.650399374969535 \times 10^{-16} \\ -0.596315604210655 \\ 0.7870714682318165 \\ -0.1578803473213 \end{pmatrix}, \quad w_2 = \begin{pmatrix} 7.085776189609232 \times 10^{-16} \\ -7.225923671895149 \times 10^{-16} \\ 0.1490848796756575 \\ -0.8058512826844632 \\ 0.5730422400206641 \end{pmatrix}, \\ w_3 = \begin{pmatrix} 0.6221573520391458 \\ -0.7312778595570425 \\ 0.01064184857958888 \\ -0.1563541004840987 \\ -0.1563541004840987 \\ 0.2315017661831681 \end{pmatrix}, \quad w_4 = \begin{pmatrix} -0.44704799767537 \\ -0.1901699845231093 \\ -0.1107453024216173 \\ -0.5888750357075574 \\ -0.6363529956501267 \end{pmatrix}, \quad w_5 = \begin{pmatrix} -1.085147452761105 \times 10^{-15} \\ -4.908431230456085 \times 10^{-16} \\ 0.120651887117552 \\ 0.6605943518791119 \\ 0.7409846316897467 \end{pmatrix}.
$$

The comparison of the FDS and the NSFDS using time series solutions for the state variables  $H(t)$  (healthy individuals), *S*(*t*) (susceptible individuals), *D*(*t*) (diabetic without complications individuals), *C*(*t*) (Diabetic with comlications individuals, and  $T(t)$  Diabetic with com lications receivin a cure individuals in figure 1). These solutions are compared using the analytical solution (green solid line), the standard finite difference scheme (FDS, blue dashed line), and the non-standard finite difference scheme (NSFDS, red dashed line with arrows).

For  $H(t)$ , the FDS differs from the analytical solution significantly after  $t = 2$ , while the NSFDS stays close to the analytical solution. Similarly, for  $S(t)$ , the FD method shows more prominent deviations after  $t = 4$ , whereas the NSFDS remains closer to the analytical solution throughout the time domain. In the case of  $D(t)$ , the FDS shows considerable differences in the decreasing phase after  $t = 2$ , while the NSFDS method accurately follows both the peak and the subsequent decrease. For  $C(t)$ , the FDS outgrows significantly during the decreasing phase after  $t = 3$ , whereas the NSFDS shows a closer agreement with the analytical solution. Finally, for  $T(t)$ , the FDS strays more as time progresses, particularly after  $t = 4$ , while the NSFDS remains close to the analytical solution throughout. Overall, the NSFDS method consistently follows the analytical solution more closely across all state variables, exhibiting better accuracy and stability, whereas the FD method exhibits larger differences, particularly as time progresses, indicating less accurate approximation.



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Figure 1. Comparison between FDS and NSFDS using time series for  $\theta = 0.913$ ,  $\beta = 0.037$ ,  $\tau = 0.035$ ,  $\mu = 0.51$ ,  $\sigma = 0.081$ ,  $\lambda_p = 0.512$ ,  $\alpha_p =$ 0.0231,  $\delta = 0.02001$ ,  $\omega = 0.079$ , and  $\gamma_p = 0.081$ 

#### **6.1** *Error magnitude*

In this section, the error between the FDS and the NSFDS is compared by using the heat maps. Both the schemes, side by side for each state variable are compared. Each heat map illustrates the error magnitude over time for the approximations. The time range is from 0 to 10, with the *y*-axis labeled "Index" representing the different time steps within this interval. The error magnitude is depicted using a color scale, ranging from 10*−*<sup>10</sup> to 10*−*<sup>2</sup> . The left panels in all the figures (Figure 2 to Figure 6) compare the error between the FDS and the analytical solution, the right panels compare the NSFDS and the analytical solution.

In the FDS heat map (left), the error is uniformly represented in a green hue across all time steps, indicating a consistent error magnitude closer to 10*−*<sup>4</sup> throughout the entire time interval. This suggests that the FDS maintains a relatively stable error but at a higher magnitude compared to the NSFDS. The NSFDS heat map (right) shows a predominantly blue to purple color, corresponding to lower error magnitudes closer to 10*−*<sup>8</sup> to 10*−*<sup>10</sup> for most time steps. This shows that the FDS maintains a uniform but higher error throughout the simulation, the NSFDS achieves significantly lower errors.



**Figure 2.** Heatmap for *H*(*t*) for error comparison between FDS and NSFDS

In the case of *S*(*t*), the FDS (left panel) exhibits errors up to 10*−*<sup>4</sup> , uniformly distributed across the time steps, indicating a consistent error magnitude throughout the simulation. On the other hand, the NSFDS (right panel) demonstrates significantly lower errors, ranging between 10*−*<sup>6</sup> and 10*−*<sup>8</sup> . This suggests that while the FDS maintains a steady but higher error, the NSFDS performs more accurately across most of the time interval.



**Figure 3.** Heatmap for *S*(*t*) for error comparison between FDS and NSFDS

In the case of  $D(t)$ , the error heat maps present a smooth pattern. For the FDS (left panel), the error remains relatively constant throughout the majority of the simulation time but shows an increase towards the upper values (closer to the end of the time interval). The error magnitude is around 10*−*<sup>4</sup> with a slight increase to 10*−*<sup>2</sup> near the end. On the other hand, the NSFDS (right panel) shows a more stable pattern where the error is much less, with values around 10*−*<sup>7</sup> to 10*−*<sup>6</sup> . This further establishes the better accuracy of the NSFDS in approximating the analytical solution for  $D(t)$ .



**Figure 4.** Heatmap for *D*(*t*) for error comparison between FDS and NSFDS

In the case of  $C(t)$ , the FDS (left panel) shows a uniform error distribution, with the magnitude of the error generally ranging from 10*−*<sup>5</sup> to 10*−*<sup>4</sup> , and highest around 10*−*<sup>2</sup> . The uniform green-yellow color suggests that the error remains fairly same throughout the simulation period, with some higher error accumulation towards the upper values for the later time steps. Oppositely, the NSFDS (right panel) displays a significantly better performance, with errors hovering around 10*−*<sup>8</sup> to 10*−*<sup>6</sup> .



**Figure 5.** Heat map for *C*(*t*) for error comparison between FDS and NSFDS

Lastly, in the case of *T*(*t*), the FDS (left panel) shows a relatively uniform error magnitudes, ranging from 10*−*<sup>5</sup> to 10*−*<sup>3</sup> , such that the highest value is at 10*−*<sup>2</sup> . This is represented by the green to yellow filling, indicating a moderate error level throughout the time steps. On the other hand, the NSFDS (right panel) demonstrates significantly better performance, with errors predominantly within the range of 10*−*<sup>8</sup> to 10*−*<sup>6</sup> .



**Figure 6.** Heat map for *T*(*t*) for error comparison between FDS and NSFDS

Across all state variables, the NSFDS demonstrates consistently lower error magnitudes compared to the standard FDS. This suggests that the NSFDS offers a more accurate and reliable approximation of the analytical solution over the given time period.

#### **6.2** *Convergence rate analysis*

In figures 7-11, the convergence rate analysis based on the step sizes and the magnitude of the error for all the state variables are discussed. The convergence rate plots display the error magnitudes for the state variable approximations using both the FDS and NSFDS for step sizes  $h = 10^{-3}$ ,  $10^{-2}$ , and  $10^{-1}$ .



**Figure 7.** Convergence rate of *H*(*t*) for error comparison between FDS and NSFDS at different step sizes



**Figure 8.** Convergence rate of *S*(*t*) for error comparison between FDS and NSFDS at different step sizes



**Figure 9.** Convergence rate of *D*(*t*) for error comparison between FDS and NSFDS at different step sizes



**Figure 10.** Convergence rate of  $C(t)$  for error comparison between FDS and NSFDS at different step sizes



**Figure 11.** Convergence rate of  $T(t)$  for error comparison between FDS and NSFDS at different step sizes

The analysis of the convergence rates reveals a uniform pattern when comparing the FDS with the NSFDS. For *H*(*t*), the FDS starts with an error of approximately 10*−*<sup>4</sup> at the smallest step size *h* = 10*−*<sup>3</sup> and slowly decreases to about <sup>6</sup>*×*10*−*<sup>5</sup> as the step size increases to *h* = 10*−*<sup>1</sup> . On the other hand, the NSFDS demonstrates a superior performance by maintaining a much lower and stable error around 10*−*<sup>6</sup> throughout all step sizes. A similar behaviour is observed for *S*(*t*), where the FDS shows a stepwise reduction in error from approximately  $10^{-4}$  at  $h = 10^{-3}$  to  $4 \times 10^{-5}$  at  $h = 10^{-1}$ , while the NSFDS uniformly holds the error around 10*−*<sup>6</sup> , signifying its accuracy. For *D*(*t*) and *C*(*t*), the FDS begins with an error close to 10*−*<sup>3</sup> and reduces to around 10*−*<sup>4</sup> as the step size increases, but it still remains considerably higher compared to the NSFDS, which maintains an error level at approximately 10*−*<sup>5</sup> across all step sizes. This uniform low error indicates the NSFDS's ability to closely follow the analytical solution, regardless of the step size. Lastly, for  $T(t)$ , the FDS shows a similar error reduction behaviour from 10*−*<sup>2</sup> to 10*−*<sup>3</sup> as the step size increases, while the NSFDS maintains an even more stable error close to 10*−*<sup>5</sup> across all tested step sizes. This detailed analysis shows the NSFDS's better accuracy and stability, making it a more reliable computational technique for accurately solving the fractional diabetes mellitus system of differential equations with varying step sizes. Overall, the NSFDS exhibits a significantly lower and more stable error magnitude across all state variables and step sizes.

### **7. Conclusion**

This paper applies the NSFDS to numerically solve the system of linear fractional differential equations that model diabetes mellitus. This scheme is found to be stable and more accurate than the FDS up to step size *h* = 10*−*<sup>3</sup> . The comparison of the FDS and the NSFDS using time series solutions for the state variables *H*(*t*) (healthy individuals),  $S(t)$  (susceptible individuals),  $D(t)$  (diabetic without complications individuals),  $C(t)$  (Diabetic with complications individuals) and  $T(t)$  (Diabetic with complications receiving a cure individuals) are presented. These solutions are compared using the analytical solution (green solid line), the standard finite difference scheme (FDS, blue dashed line) and NSFDS, red dashed line with arrows. The NSFDS method consistently follows the analytical solution more closely across all state variables, indicating better accuracy and stability, whereas the FDS exhibits larger differences, particularly as time progresses, indicating less accurate approximation. This demonstrates the better performance of the NSFDS method in approximating the analytical solution of the fractional diabetic model. The heat maps compare the errors between the analytical solution and two numerical methods: the standard finite difference FDS and the NSFDS. Each heat map illustrates the error magnitude over time for the state variables. Across all state variables, the NSFDS demonstrates consistently lower error magnitudes compared to the standard FDS. This suggests that the NSFDS offers a more accurate and reliable approximation of the analytical solution over the given time period. The heat maps effectively visualize these differences, with the NSFDS providing better performance in reducing errors for the fractional diabetic model. The convergence rate analysis discussed here, is based on the step sizes and the magnitude of the error for all the state variables. The convergence rate plots display the error magnitudes using both the FDS and NSFDS for step sizes  $h = 10^{-3}$ ,  $10^{-2}$ , and 10*−*<sup>1</sup> .

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# **Confilict of interest**

The authors declare no competing financial interest.

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